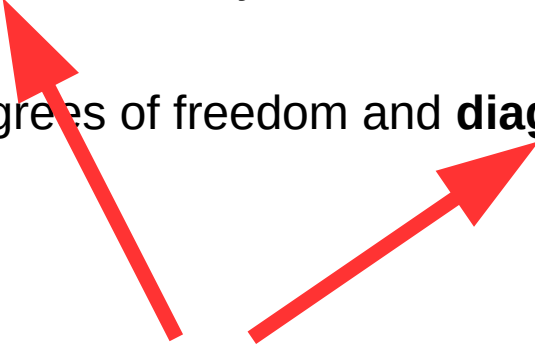


Hydrostatic modelling

- **What** is the hydrostatic approximation ? When is it **valid** ?
- What are the **degrees of freedom** of a hydrostatic model ?
- How do we **prognose** the degrees of freedom and **diagnose** other quantities ?

answer depends on the choice of *vertical coordinate*



Hydrostatic approximation : basic idea and order-of-magnitude arguments

Vertical momentum balance :

$$\frac{dw}{dt} = -g - \frac{1}{\rho} \frac{\partial p}{\partial z}$$

If vertical acceleration is small :

$$\frac{dw}{dt} \ll g, \quad \frac{1}{\rho} \frac{\partial p}{\partial z} \quad \frac{\partial p}{\partial z} = -\rho g$$

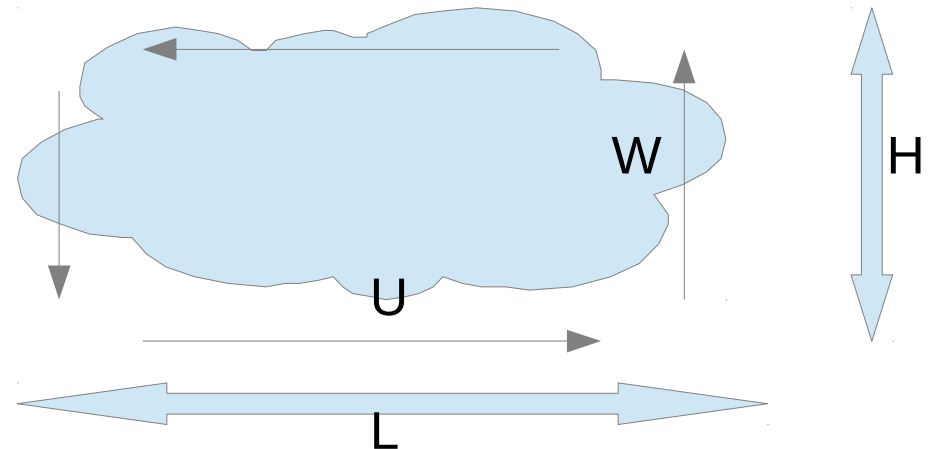
How good is this approximation ?
what we neglect << what is left ?

$$\frac{d}{dt} \sim \frac{W}{H} \sim \frac{U}{L} \quad \frac{dw}{dt} \sim \frac{W^2}{H}$$

$$\frac{\partial p}{\partial x} \sim \rho \frac{du}{dt} \sim \rho \frac{U}{L} U$$

$$\frac{\partial p}{\partial z} \sim \frac{L}{H} \frac{\partial p}{\partial x} \sim \frac{\rho U^2}{H}$$

$$\rho \frac{Dw}{Dt} / \frac{\partial p}{\partial z} \sim \left(\frac{W}{U} \right)^2 \sim \left(\frac{H}{L} \right)^2$$



Hydrostatic approximation
valid if

$$\frac{H}{L} \ll 1$$

Atmosphere : OK for **large-horizontal-scale circulation** (>30km)
KO for convection (storms), orographic flow (mountains).

Ocean : H~1km => OK down to **kilometer-scale**

Hydrostatic approximation from least action principle

$$\frac{D\mathbf{m}}{Dt} - \frac{\partial L}{\partial \mathbf{x}} - \frac{1}{\rho} \nabla \left(\rho^2 \frac{\partial L}{\partial \rho} \right) = 0 \quad \mathbf{m} = \frac{\partial L}{\partial \mathbf{u}} \quad PV = \frac{1}{\rho} \nabla_s \cdot (\nabla \times \mathbf{m})$$

$$L = \frac{\mathbf{u} \cdot \mathbf{u}}{2} + \mathbf{R} \cdot \mathbf{u} - \Phi - h(p, s) + \frac{p}{\rho}$$

$$\frac{D\mathbf{u}}{Dt} + \text{curl} \mathbf{R} \times \mathbf{u} + \nabla \Phi + \frac{1}{\rho} \nabla p = 0$$

$$E = \frac{\mathbf{u} \cdot \mathbf{u}}{2} + \Phi + h(p, s) - \frac{p}{\rho}$$

$$\mathbf{m} = \mathbf{u} + \mathbf{R}$$



Horizontal projection

$$\mathbf{u}_H = \mathbf{u} - u_z \mathbf{e}_z$$

$$L = \frac{\mathbf{u}_H \cdot \mathbf{u}_H}{2} + \mathbf{R} \cdot \mathbf{u}_H - \Phi - h(p, s) + \frac{p}{\rho}$$

$$\frac{D\mathbf{u}_H}{Dt} + \text{curl} \mathbf{R} \times \mathbf{u}_H + \nabla \Phi + \frac{1}{\rho} \nabla p = 0$$

$$E = \frac{\mathbf{u}_H \cdot \mathbf{u}_H}{2} + \Phi + h(p, s) - \frac{p}{\rho}$$

$$\mathbf{m} = \mathbf{u}_H + \mathbf{R}$$

Prognosing hydrostatic + Boussinesq

u, v, s, q



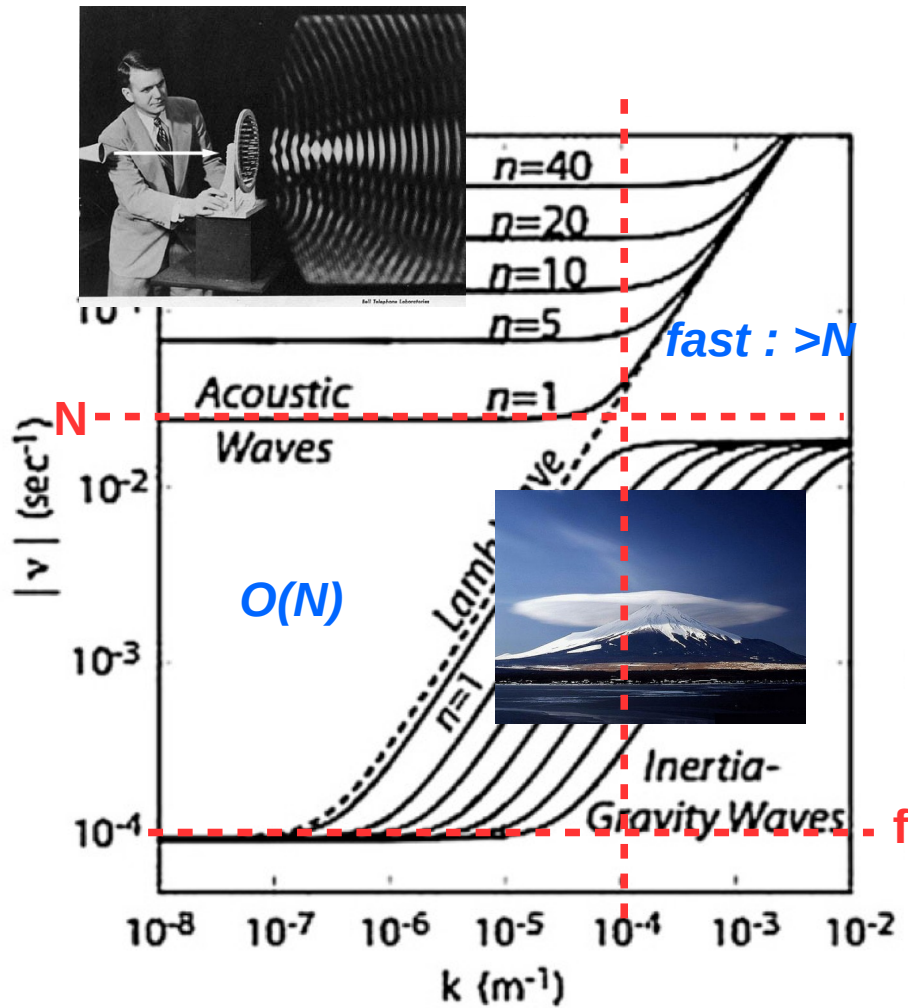
$$\begin{aligned}
 &\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} + \frac{\partial w}{\partial z} = 0 \\
 &\frac{1}{\rho_0} \frac{\partial p'}{\partial z} = gb(z, s, q) \\
 &\frac{\partial q}{\partial t} + \mathbf{u} \cdot \nabla q = 0 \\
 &\frac{\partial s}{\partial t} + \mathbf{u} \cdot \nabla s = 0 \\
 &\frac{\partial \mathbf{u}_H}{\partial t} + (\mathbf{u} \cdot \nabla) \mathbf{u}_H + f \mathbf{e}_z \times \mathbf{u}_H + \frac{1}{\rho_0} \nabla_H p' = 0
 \end{aligned}$$



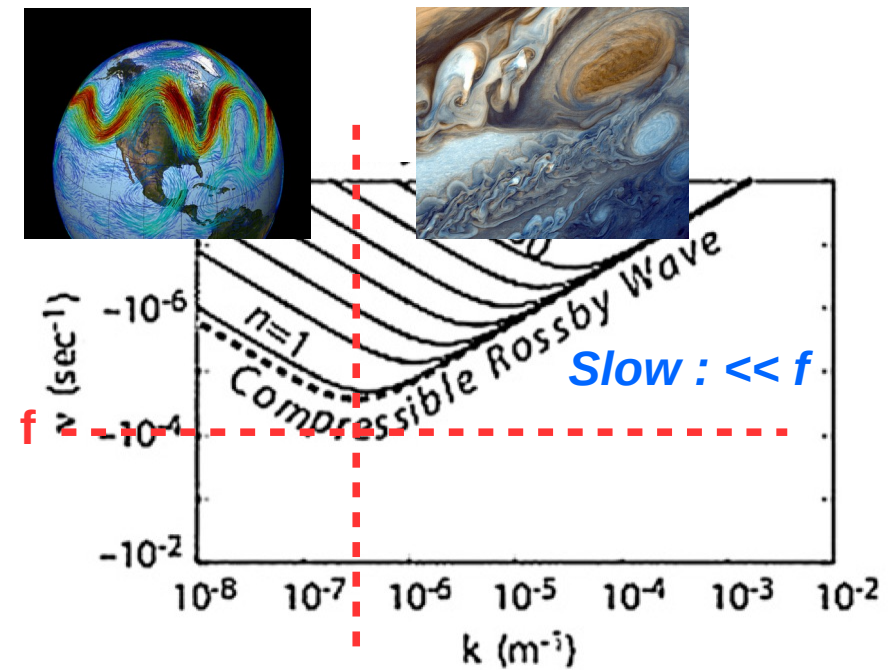
$\frac{\partial u}{\partial t}, \frac{\partial v}{\partial t}, \frac{\partial s}{\partial t}, \frac{\partial q}{\partial t}$

Classes of motion in a gravity-dominated, rotating, compressible flow :

Waves in an isothermal atmosphere at rest



f-plane

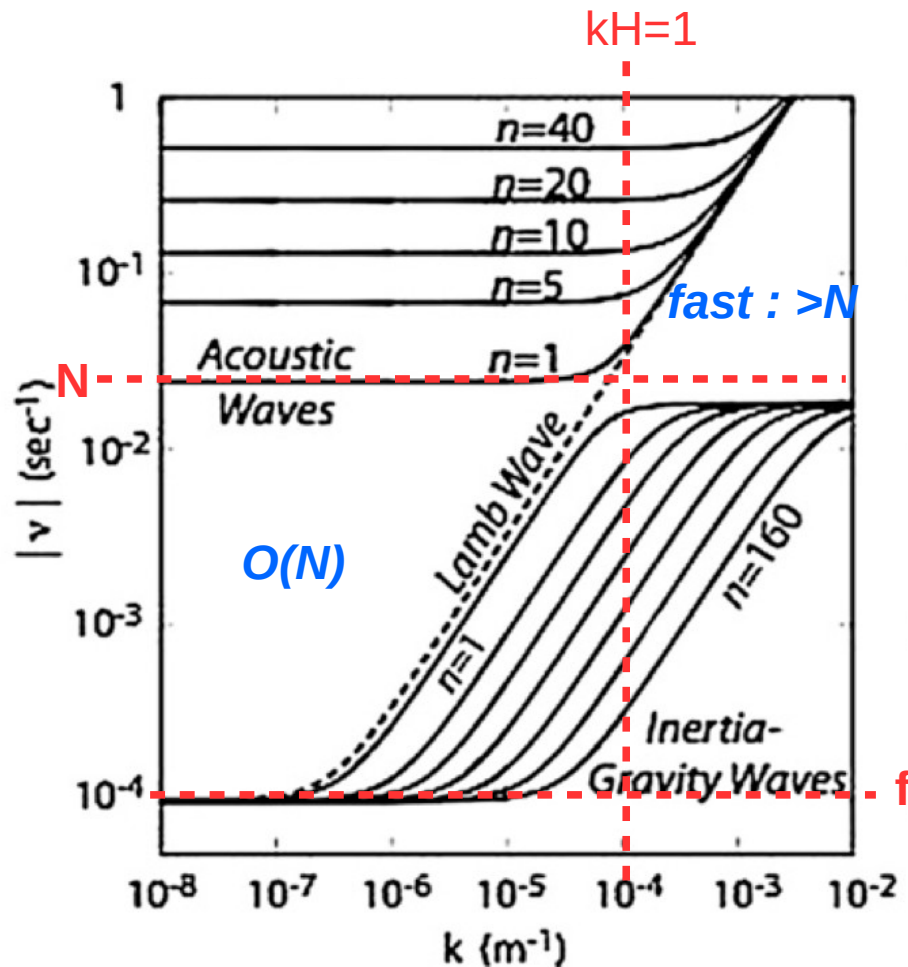


beta-plane

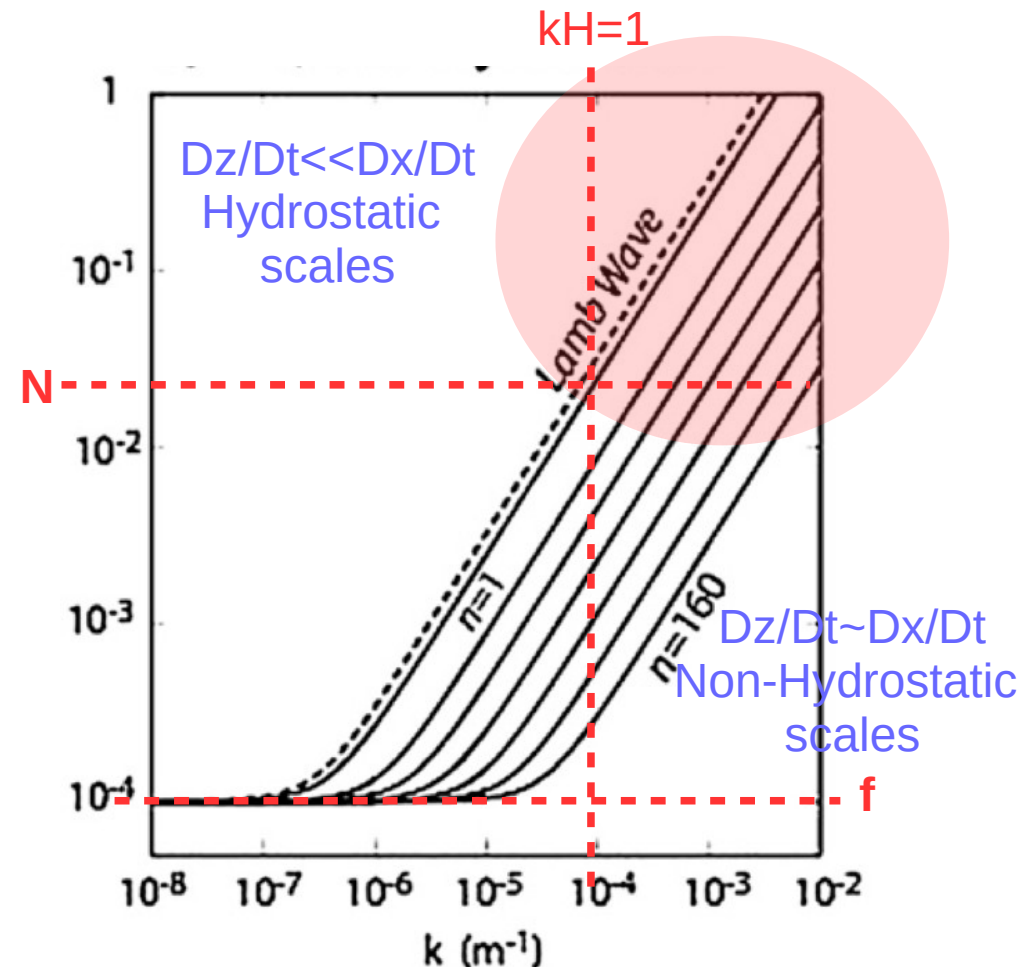
Waves in an isothermal atmosphere at rest

Traditional f-plane approximation

« Exact »



Hydrostatic



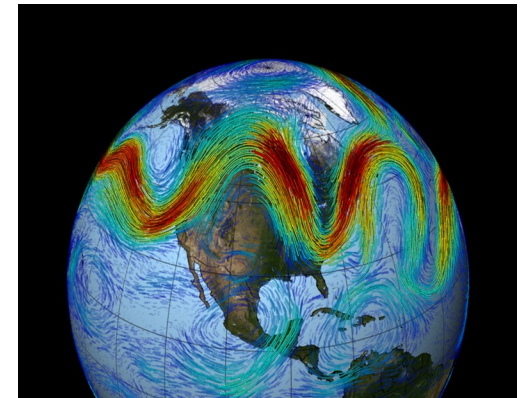
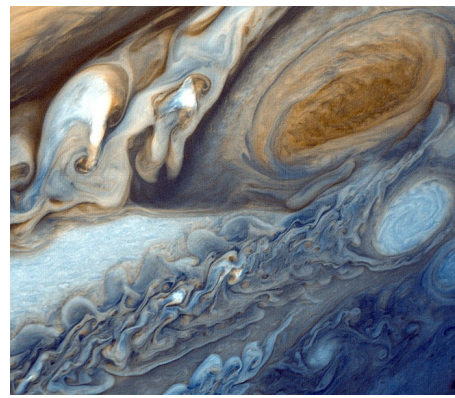
$$\frac{dw}{dt} / \rho \frac{\partial p}{\partial z} \sim \left(\frac{W}{U} \right)^2 \sim \left(\frac{H}{L} \right)^2$$

Characteristic scales

- *Velocity* : Sound $c \sim 340\text{m/s}$ Wind $U \sim 30\text{m/s}$
- *Time* : Buoyancy oscillations $N \sim g/c \sim 10^{-2} \text{ s}^{-1}$ Coriolis $f \sim 10^{-4} \text{ s}^{-1}$
- *Length* : **Scale height $H=c^2/g=10\text{km}$** **Rossby radius : $R=c/f \sim 1000 \text{ km}$**

Mach number : $M=U/c \ll 1$

Scale separation : $f/N \sim H/R \ll 1$



small-scale

scale height

mesoscale

synoptic

planetary

1 km

10 km

100 km

1000 km

10000 km

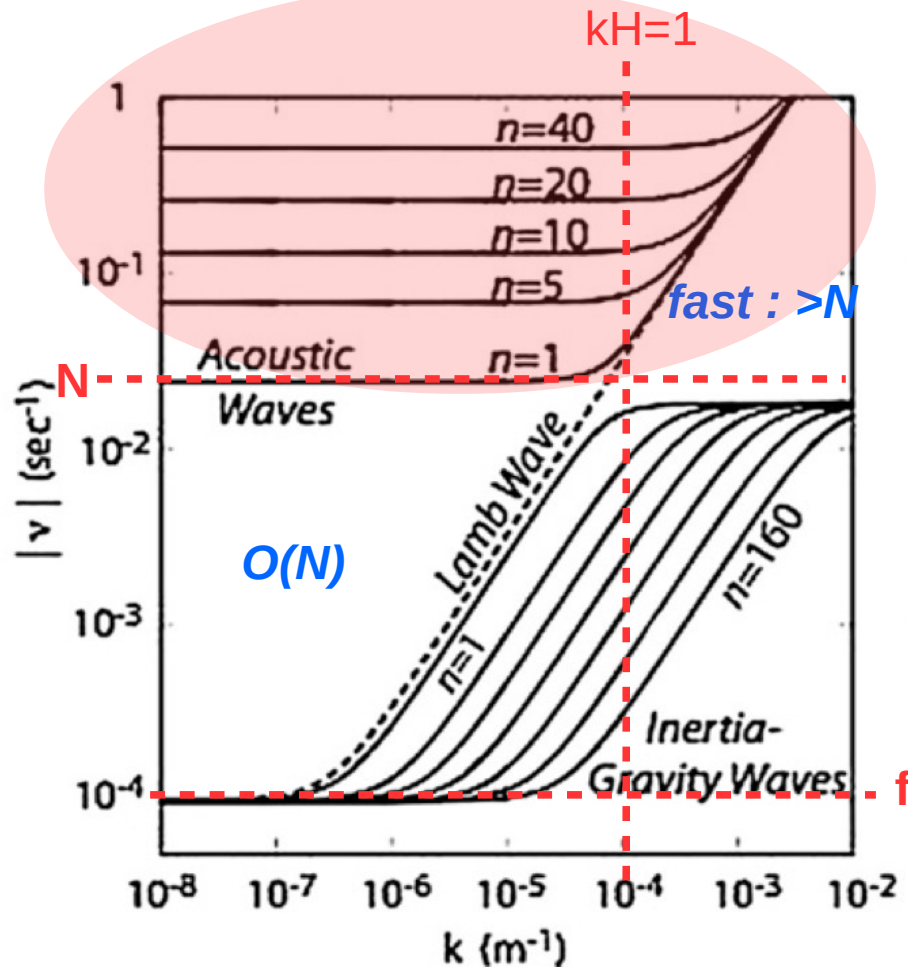
non-hydrostatic

hydrostatic

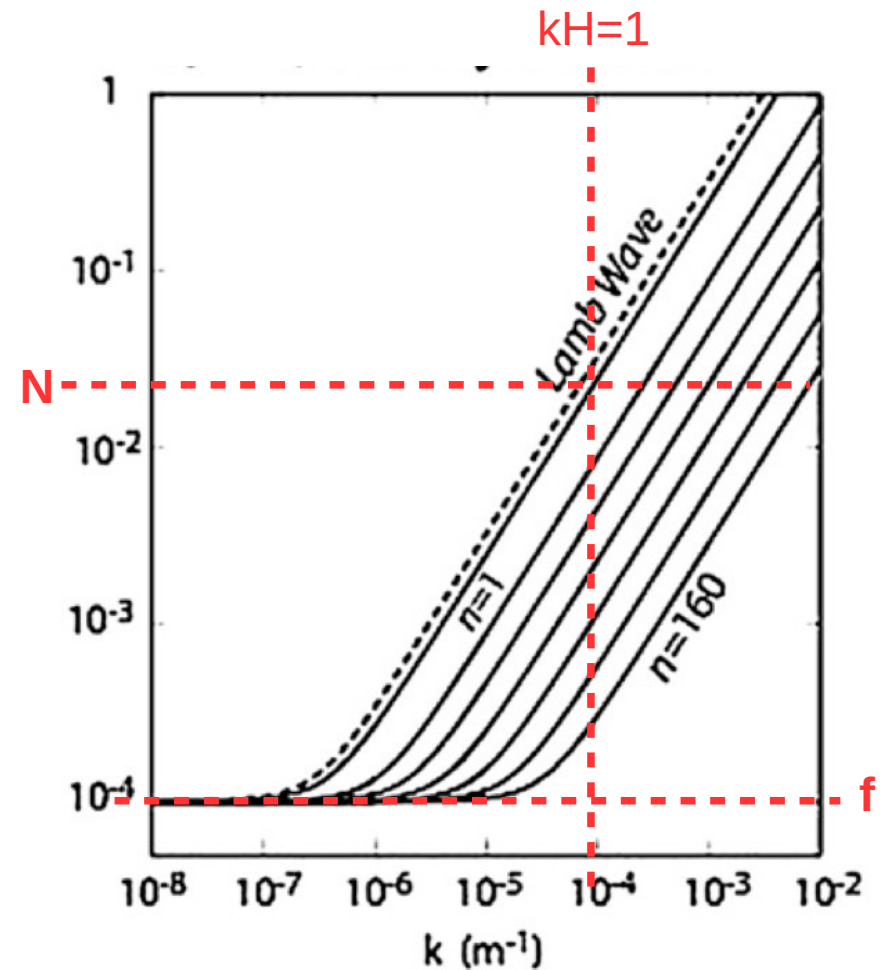
Waves in an isothermal atmosphere at rest

Traditional f -plane approximation

Euler



Hydrostatic



Acoustic waves suppressed : « sound-proof » approximation

Spherical geoid

Shallow-atmosphere +
traditional

(Quasi-)Hydrostatic

Boussinesq / Anelastic /

Pseudo-incompressible

Boussinesq

Anelastic
(*Ogura & Phillips*)

Pseudo-
incompressible
(*Durran ; Klein &
Pauluis*)

Acoustic
Lamb
Inertia-gravity
Rossby

Compressible
Euler

Spherical-geoid
Euler

Traditional shallow-
atmosphere
(*Phillips, 1966*)

Acoustic
Lamb
Inertia-gravity
Rossby

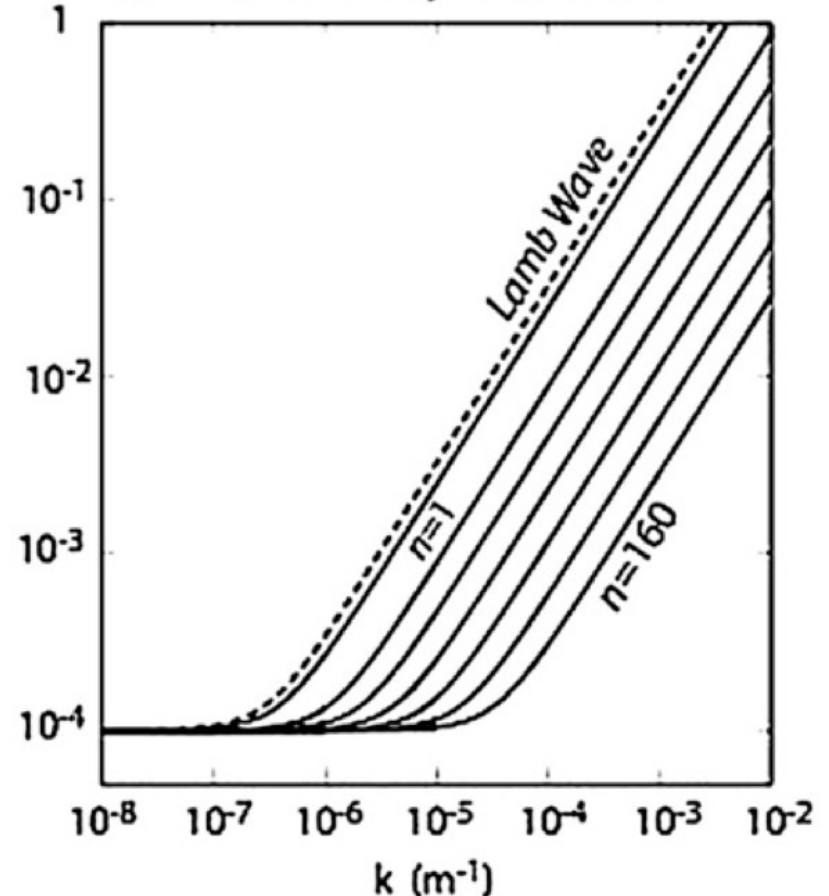
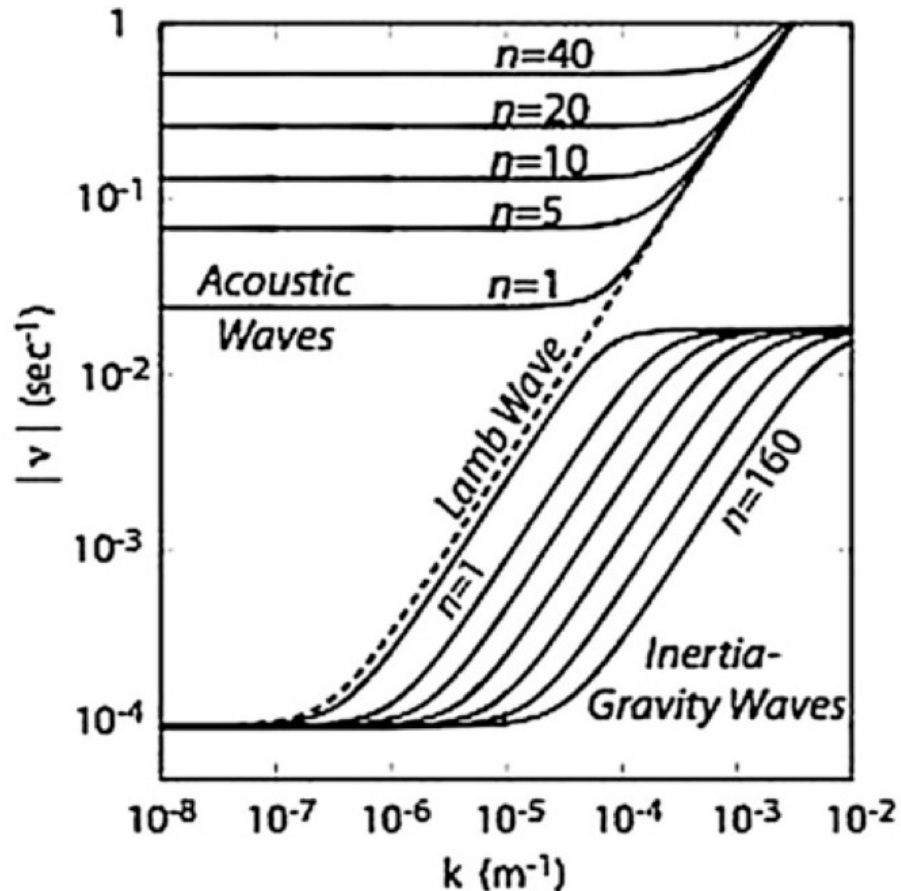
Quasi-hydrostatic
(*White & Wood, 2012*)

Spherical-geoid
Quasi-hydrostatic
(*White & Wood, 1995*)

Primitive equations
(*Richardson, 1922*)

Acoustic
Lamb
Inertia-gravity
Rossby

Sound-proofing and degrees of freedom



- Euler equations : 5 prognostic fields (u, v, w , density, entropy)
 $\Rightarrow$ linearized equations lead to a 5×5 determinant
 $\Rightarrow$ degree-5 algebraic equation for $\omega(k)$
 $\Rightarrow$ 5 possible values (roots) for $\omega(k)$:
 2 acoustic, 2 inertia-gravity, 1 Rossby

Hydrostatic :

- only 3 roots : inertia-gravity, Rossby
 $\Rightarrow$ only 3 independent prognostic fields !
- Hydrostatic balance acts as a *constraint* that reduces the numbers of degrees of freedom
- Which fields have become diagnostic ?
- How do I diagnose other fields ?

Compressible hydrostatic dynamics :

are we ready to prognose ?

$$\partial_z p(\rho, s) + \rho g = 0$$

$$\partial_t \rho + \partial_x(\rho u) + \partial_y(\rho v) + \partial_z(\rho w) = 0$$

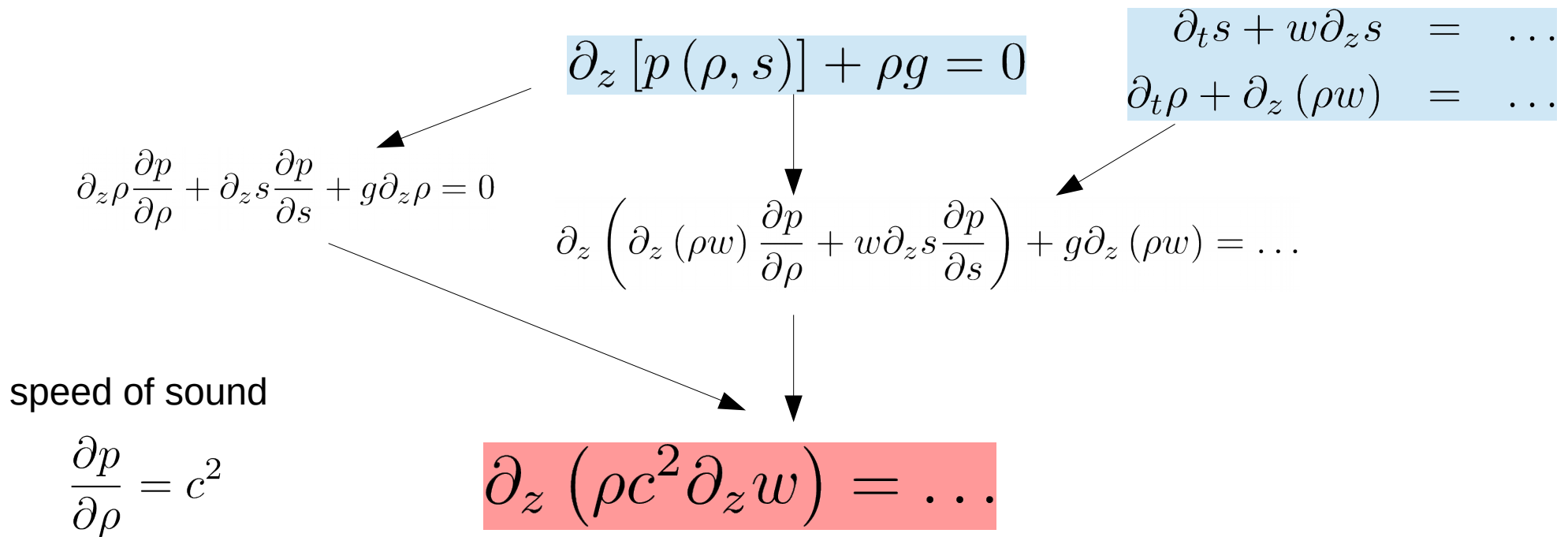
$$\partial_t s + u \partial_x s + v \partial_y s + w \partial_z s = 0$$

$$\partial_t u + u \partial_x u + v \partial_y u + w \partial_z u - f v + \frac{1}{\rho} \partial_x p = 0$$

$$\partial_t v + u \partial_x v + v \partial_y v + w \partial_z v + f u + \frac{1}{\rho} \partial_y p = 0$$

- 4 prognostic equations for 3 degrees of freedom ??
- Hydrostatic balance yields density given entropy
- How do we diagnose vertical velocity w ?
- There must be an additional diagnostic relationship = *hidden constraint*
- In order to obtain hidden constraints, time-differentiate known constraints (here : hydrostatic balance)

Compressible hydrostatic dynamics : prognostic and diagnostic degrees of freedom



Time-differentiating hydrostatic balance
=> diagnostic equation for vertical velocity

- First obtained by Richardson (1922) but using pressure as a prognostic variable
- Ooyama (1990) : prognosing entropy => « neater form » of Richardson's equation
- Dubos & Tort (2014) : General form for quasi-hydrostatic systems in curvilinear coordinates

Compressible hydrostatic dynamics :

are we ready to prognose ?

$$\partial_z p(\rho, s) + \rho g = 0$$

$$\partial_z (\rho c^2 \partial_z w) = \dots$$

$$\partial_t s + u \partial_x s + v \partial_y s + w \partial_z s = 0$$

$$\partial_t u + u \partial_x u + v \partial_y u + w \partial_z u - f v + \frac{1}{\rho} \partial_x p = 0$$

$$\partial_t v + u \partial_x v + v \partial_y v + w \partial_z v + f u + \frac{1}{\rho} \partial_y p = 0$$

- Hydrostatic balance yields density given entropy
- Richardson's equation yields vertical velocity
- 3 prognostic equations for 3 degrees of freedom

- Typical compressible hydrostatic models do not work this way
- The reason can be found by having a closer look at *hydrostatic adjustment*

Hydrostatic adjustment

or : *how nature imposes hydrostatic balance*

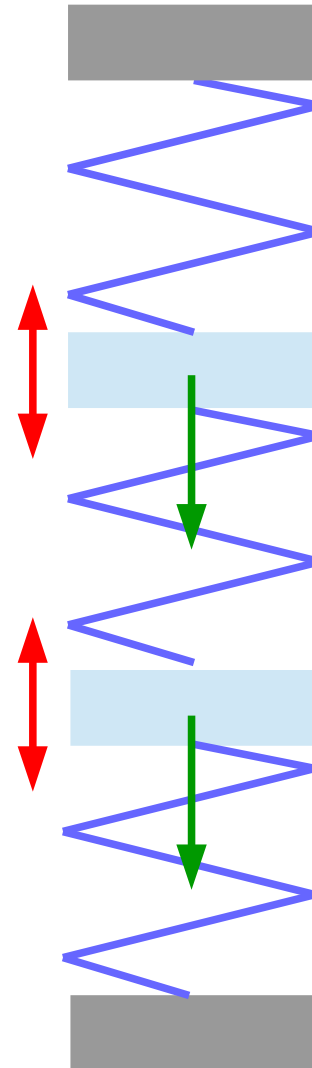
Consider

a horizontally homogeneous atmosphere
initial profile $\rho(z)$, $s(z)$ not hydrostatically balanced
what happens ?

Mechanical analogy :

a vertical stack of masses subject to gravity
and coupled by springs

- Vertical forces initially unbalanced
- Vertical acceleration, vertical displacements
- Oscillations / *acoustic waves*
- Until a balance is reached eventually
- Balanced state minimizes *mechanical energy*



Hydrostatic adjustment

or : *how nature imposes hydrostatic balance*

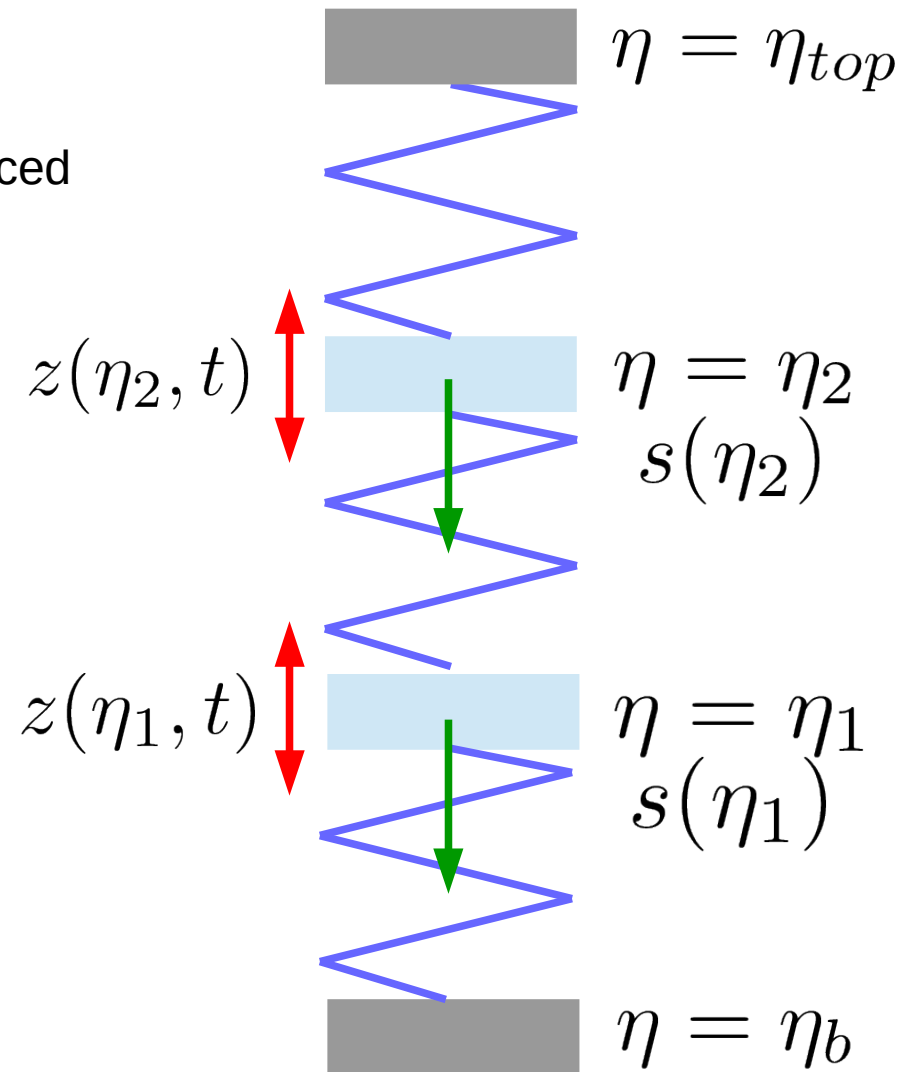
Consider

a horizontally homogeneous atmosphere
initial profile $\rho(z)$, $s(z)$ not hydrostatically balanced
what happens ?

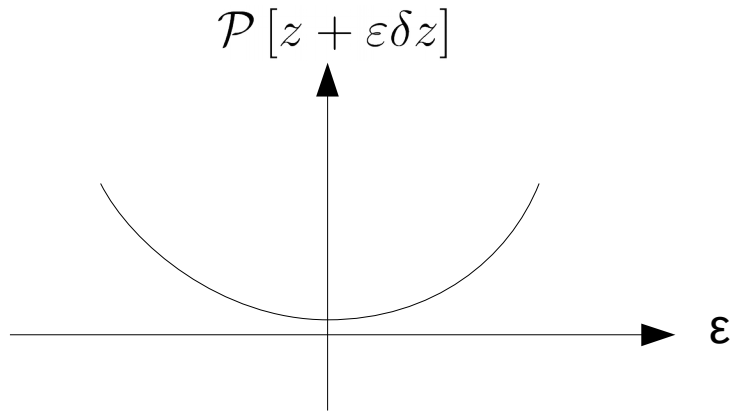
Mechanical analogy :

a vertical stack of masses subject to gravity
and coupled by springs

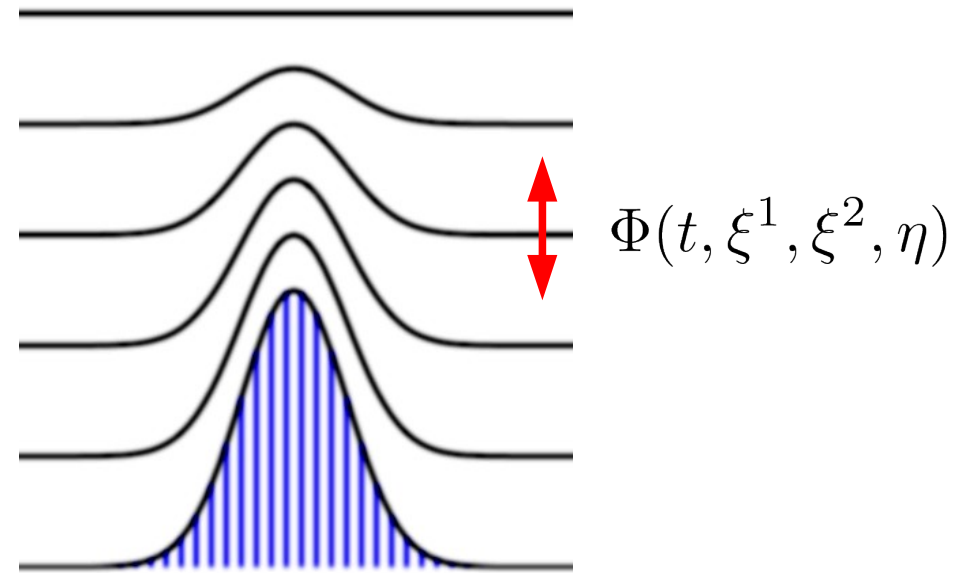
- Vertical forces initially unbalanced
- Vertical acceleration, vertical displacements
- Oscillations / *acoustic waves*
- Until a balance is reached eventually
- Balanced state minimizes *mechanical energy*



Hydrostatic adjustment and vertical coordinate



- Hydrostatic adjustment : 1D, well-posed, nonlinear elliptic problem where *the altitude z of air parcels is the unknown*
- In a hydrostatic model, a hydrostatic adjustment occurs at each time step suggests to :
 - let model layers « float » vertically
 - let altitude z be a *field* rather than coordinate
 - Non-Eulerian vertical coordinate



$$\partial_t \mu + \partial_1(\mu u^1) + \partial_2(\mu u^2) + \partial_\eta(\mu \dot{\eta}) = 0$$

$$\partial_t s + u^1 \partial_1 s + u^2 \partial_2 s + \dot{\eta} \partial_\eta s = 0$$

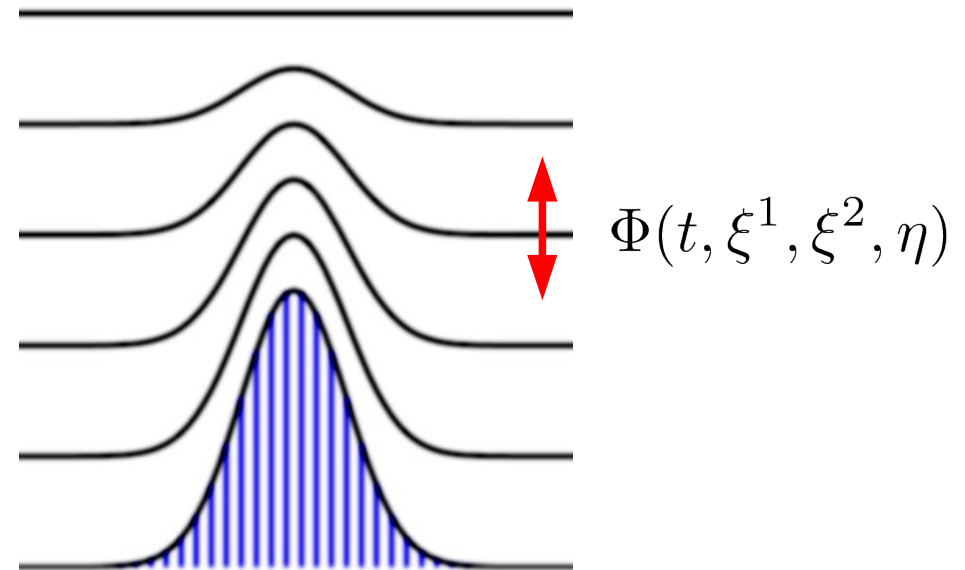
Procedure to *diagnose* $\dot{\eta}$ depends on specific definition of vertical coordinate. This definition is purely *kinematic*.

Generalized vertical coordinates

Isentropic vertical coordinate $s = s(\eta)$

- Entropy becomes *diagnostic*
- Vertical « velocity » given by heating
- Ground usually not an isentrope ...
- Ocean surface usually not an isopycnal...

*Purely isentropic/isopycnal
coordinate not used in practice*

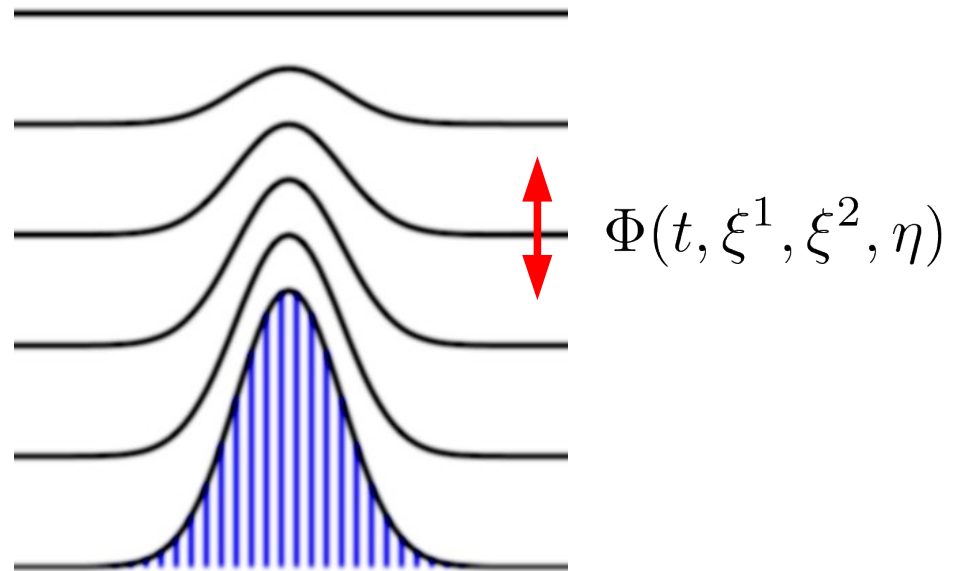


$$\partial_t \mu + \partial_1(\mu u^1) + \partial_2(\mu u^2) = 0$$
$$s = s(\eta)$$

Generalized vertical coordinates

Lagrangian vertical coordinate $\dot{\eta} = 0$

- simple
- initial altitude of layers can be chosen arbitrarily
- layers do not exchange mass, entropy
=> they follow the flow
=> layers tend to *fold* and *cross*



$$\partial_t \mu + \partial_1(\mu u^1) + \partial_2(\mu u^2) = 0$$

$$\partial_t s + u^1 \partial_1 s + u^2 \partial_2 s = 0$$

solution :

- exploit arbitrariness of altitude
- *vertical remap* before layers fold/cross

Generalized vertical coordinates

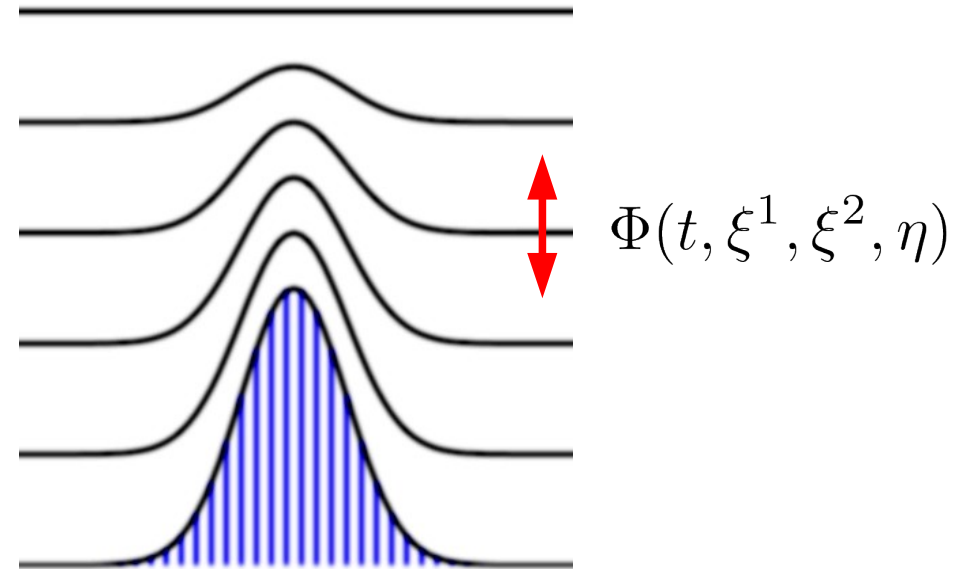
Hybrid mass-based coordinate

- each layer contains a prescribed fraction of the **total mass of the column** + fixed amount

$$\mathbf{M}(t, \xi^1, \xi^2) \equiv \int \mu d\eta$$

$$\mu = a(\eta)\mathbf{M} + b(\eta)$$

- Layers exchange mass in order to respect this prescription
- There is always some mass in each layer, so layers never fold/cross



- Diagnose μ from total column mass $\mathbf{M}$
- Prognose $\mathbf{M}$
- Diagnose $d\mu/dt$
- Diagnose $\eta_{\dot{}}$

$$\mu = a(\eta)\mathbf{M} + b(\eta)$$

$$\partial_t \mathbf{M} + \partial_1 \int \mu u^1 d\eta + \partial_2 \int \mu u^2 d\eta = 0$$

$$\partial_t \mu = A(\eta) \partial_t \mathbf{M}$$

$$\partial_t \mu + \partial_1 (\mu u^1) + \partial_2 (\mu u^2) + \partial_\eta (\mu \dot{\eta}) = 0$$

Generalized vertical coordinates

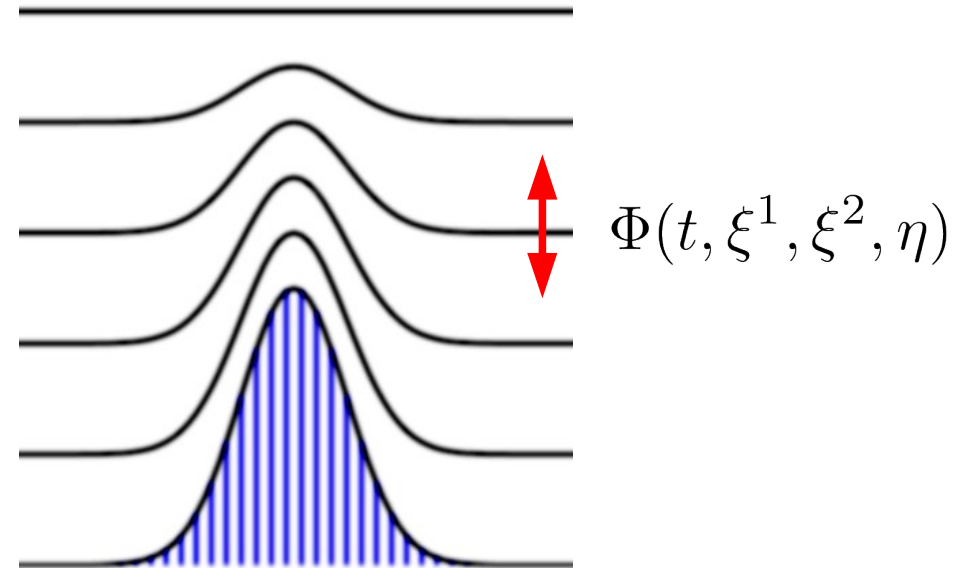
Hybrid pressure coordinate

Special form of mass-based, if we are solving **traditional hydrostatic** equations

$$\partial_{\eta} p = -\mu g$$

$$M(t, \xi^1, \xi^2) \equiv \int \mu d\eta \Rightarrow p_s = p_{top} + gM$$

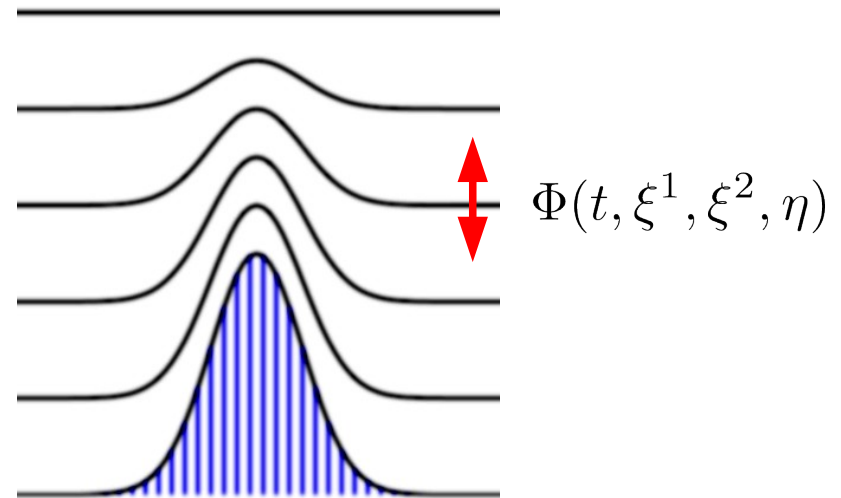
$$\mu = a(\eta)M + b(\eta) \Rightarrow p = A(\eta)p_s + B(\eta)$$



- Prognose **surface pressure** rather than M
- Diagnose pressure straightforwardly
No need to solve hydrostatic balance
- Hybrid coefficient A can be chosen close to 0
so that **upper layers are nearly isobars**

Recap : hydrostatic dynamics, generalized vertical coordinates & prognostic variables

- a **hydrostatic adjustment** occurs at each time step
- altitude z : time-dependent diagnostic field rather than coordinate
- **Non-Eulerian vertical coordinate**



Hybrid mass-based coordinate

- Diagnose pseudo-density μ from total column mass M
- Prognose M
- Diagnose $d\mu/dt$
- Diagnose $\eta_{\dot{}}$
- Prognose entropy
- **Hydrostatic adjustment** => geopotential
- Prognose momentum

Lagrangian coordinate

- Prognose pseudo-density μ
- Prognose entropy
- If needed, vertical remap
- **Hydrostatic adjustment** => geopotential
- Prognose momentum

kinematics

dynamics

References

- Phillips (1957) J. Meteorology **14**:184-185
- Kasahara (1974) Mon. Wea. Rev. **102**: 509-522
- Ooyama (1990) J. Atmos. Sci **47**(21) : 2580-2593
- Laprise (1992) Mon. Wea. Rev. **120** : 197-207
- Dubos & Tort (2014) Mon. Wea. Rev. **142**(10) : 3860-3880